

1 **Physical controls of dynamics of methane venting from a shallow seep area west of**  
2 **Svalbard**

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23 **Eddies; Ocean model; Svalbard shelf**

24 **Highlights**

- 25 • Gas seepage intensity and lateral water mass movements are key controls of water column  
26 methane content.
- 27 • Vertical methane transport is limited irrespective of stratification.
- 28 • Eddies play a key role in horizontal advection and dispersion of dissolved methane.

29

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## 31 **Abstract**

32 We investigate methane seepage on the shallow shelf west of Svalbard during three  
33 consecutive years, using discrete sampling of the water column, echosounder-based gas flux  
34 estimates, water mass properties, and numerical dispersion modelling. The results reveal three  
35 distinct hydrographic conditions in spring and summer, showing that the methane content in the  
36 water column is controlled by a combination of free gas seepage intensity and lateral water mass  
37 movements, which disperse and displace dissolved methane horizontally away from the seeps.  
38 Horizontal dispersion and displacement of dissolved methane are promoted by eddies originating  
39 from the West Spitsbergen Current and passing over the shallow shelf, a process that is more  
40 intense in winter and spring than in the summer season. Most of the methane injected from  
41 seafloor seeps resides in the bottom layer even when the water column is well mixed, implying  
42 that the controlling effect of water column stratification on vertical methane transport is small.  
43 Only small concentrations of methane are found in surface waters, and thus the escape of  
44 methane into the atmosphere above the site of seepage is also small. The magnitude of the sea to  
45 air methane flux is controlled by wind speed, rather than by the concentration of dissolved  
46 methane in the surface ocean.

## 47 **1 Introduction**

48 The Arctic Ocean holds vast reservoirs of the potent greenhouse gas methane in the form  
49 of free and dissolved gas (Lammers et al., 1995; Damm et al., 2005), gas entrapped in subsea  
50 permafrost (Shakhova et al., 2010), and gas hydrates in sediments (Hester and Brewer, 2009;  
51 Westbrook et al., 2009; Berndt et al., 2014). In particular, gas that is bound in hydrates may be  
52 released as a result of temperature induced gas hydrate destabilization (Kretschmer et al., 2015;  
53 James et al., 2016), which makes the warming Arctic Ocean a potential hot spot of future  
54 methane emission (Shakhova et al., 2010; Kort et al., 2012; Parmentier et al., 2015). Methane  
55 release from the seafloor has been documented from numerous areas along the Arctic Ocean  
56 continental margin: the West Spitsbergen continental margin and shelf (Knies et al., 2004; Damm  
57 et al., 2005; Westbrook et al., 2009; Sahling et al., 2014; Smith et al., 2014; Graves et al., 2015;  
58 Mau et al., 2017), the Barents Sea (Lammers et al., 1995; Serov et al., 2017; Andreassen et al.,  
59 2017), the Kara Sea shelf (Portnov et al., 2013; Serov et al., 2015), the East Siberian Shelf  
60 (Shakhova et al., 2010, 2013), and the Beaufort Sea (Kvenvolden et al., 1993; Paull et al., 2007).

61 Methane release from the West Spitsbergen margin particularly has been ongoing for several  
62 millennia and is, at least partly, temperature controlled (Berndt et al., 2014).

63 Indeed, Arctic air temperatures are increasing twice as fast as the global average because  
64 of Arctic amplification (Graversen et al., 2008; Serreze and Francis, 2006; Pachauri et al., 2014).  
65 The annual average Arctic air temperature is now 3.5°C warmer compared to the beginning of  
66 20th century (Soreide et al., 2016). As a result, expanding areas of ice-free Arctic Ocean waters  
67 are being exposed to solar radiation and elevated air temperatures. Combined with an increase of  
68 heat input from adjacent ocean basins, e.g. warmer than usual Atlantic Water (AW) propagating  
69 deeper into the Arctic Ocean (Polyakov et al., 2004; 2007; 2010), this results in a present day  
70 Arctic Ocean sea surface temperature, which is 5°C warmer than the 1982-2010 average for the  
71 Barents and Chukchi seas and around Greenland (Soreide et al., 2016). The effect of increasing  
72 temperature in the future Arctic may therefore become more important for Arctic seafloor  
73 methane liberation (Westbrook et al., 2009; Ferré et al., 2012; Marín-Moreno et al., 2015).

74 Several processes determine the fate of methane released into the water column from  
75 sediments and, most importantly, its release to the atmosphere. Methane contained in bubbles  
76 emanating from the seafloor dissolves in seawater and can be rapidly transported from the area by  
77 the advection of water masses (Graves et al., 2015). The upward transport of dissolved methane  
78 has been found to be limited by water column stratification (e.g. Schmale et al., 2005; Leifer et  
79 al., 2009). Studies by Myhre et al., (2016) and Gentz et al., (2014) conducted on the shallow shelf  
80 and upper continental slope off Svalbard, west of Prins Karls Forland (PKF) revealed waters  
81 enriched with dissolved methane below the pycnocline. However, the methane concentrations  
82 above the pycnocline were generally in equilibrium with the atmospheric mixing ratio. This  
83 suggested that the pycnocline may act as a physical barrier, preventing dissolved methane from  
84 entering the well mixed upper layer of the water column and thus also the atmosphere, instead  
85 trapping methane in the lower sphere of the water column. The open Arctic Ocean is stratified  
86 throughout the year (Rudels et al., 1994). In shallower areas, however, the stratification of the  
87 entire water column is subject to an annual cycle and a seasonal erosion of the pycnocline e.g.  
88 through winter time convection or wind induced mixing (Cottier et al., 2010). If controlled by  
89 stratification, the escape of methane to the atmosphere would also follow this seasonality. In  
90 other words, the potential for methane to be liberated to the atmosphere from these areas is higher  
91 when there is no stratification during stormy seasons (von Deimling et al., 2011).

92 Another important process determining the fate of methane in the water column is its  
93 removal by aerobic methane oxidation (MOx), mediated by aerobic methanotrophic bacteria  
94 (Hanson & Hanson, 1996; Reeburg, 2007; Steinle et al., 2015). Methane removal from deep  
95 water sources through MOx is more efficient than that from shallow sources, because the distance  
96 between methane liberation from the seafloor and potential methane evasion to the atmosphere is  
97 greater and methanotrophs in the water column have more time for methane consumption (Steinle  
98 et al., 2015; Graves et al., 2015; James et al., 2016). For example, in the deep Gulf of Mexico  
99 (~1500 meters water depth), most of the methane discharged following the Deepwater Horizon  
100 oil spill was consumed by water column methanotrophs (Kessler et al., 2011), while most of  
101 methane seeping from the shallow seafloor on the East-Siberian Shelf (~50 m water depth) was  
102 liberated to the atmosphere, especially during storm-induced mixing events (Shakhova et al.,  
103 2013).

104 Cold seeps methane is also used by a variety of free-living and symbiotic chemosynthetic  
105 benthic organisms (e.g. Knittel & Boetius 2009; Levin et al., 2016, Åström et al., 2018). A key  
106 process at cold seeps is methane oxidation (anaerobic or aerobic) producing CO<sub>2</sub>/DIC (dissolved  
107 inorganic carbon) and hydrogen sulphide, which in turn can be used by other organisms. In  
108 addition, heterotrophic organisms graze on chemosynthetic autotrophs, thereby passing  
109 chemosynthetically fixed carbon to higher levels of the food web (Niemann et al., 2013, Åström  
110 et al., 2019). Cold seeps often support high standing stocks of biomass (Sahling et al., 2003;  
111 Boetius and Suess, 2004). However, shallow water cold seeps are characterised by a lower  
112 species richness and biomass when compared to deep water seeps, possibly because of a higher  
113 grazing pressure at shallow water systems (Sahling et al., 2003). Cold seeps share some  
114 characteristics with hydrothermal systems (Paull et al., 1984; Levin et al., 2016), i.e. both systems  
115 are fuelled by reduced compounds from the geosphere. Yet, the driving geological mechanisms  
116 leading to cold seeps and hydrothermal vents differ. Hydrothermal vents are driven by magmatic  
117 processes, while cold seeps are driven by a variety of, often sedimentary processes, such as clay  
118 mineral dewatering. In addition, hydrothermal vents are primarily enriched in sulphide and not  
119 methane as in the typical case for cold seeps.

120 Marine environments in the Arctic Ocean characterised by ongoing methane release are  
121 ideal natural laboratories for studying the effects of potentially enhanced seafloor methane  
122 venting in warming waters, and the processes that regulate the transport of this methane. In this

123 paper, we study the dynamics of methane venting from shallow gas-bearing sediments (water  
124 depth: 50-120 meters) west of PKF off the Svalbard archipelago; and the physical processes in  
125 the water column that control methane dispersion and displacement away from the seeps. We  
126 conducted hydroacoustic surveys to determine the flux of free gas (i.e. bubbled methane) from  
127 sediments, along with oceanographic surveys to determine concentrations of dissolved methane  
128 in the water column, sea-air methane fluxes, and water mass properties. Measurements were  
129 repeated in a defined study area during three consecutive years to investigate the dynamics of  
130 venting methane under varying hydrographic conditions. Model simulations place these detailed  
131 observations into the broader seasonal context, and allow a better understanding of the  
132 oceanographic processes controlling methane dynamics in the area of study.

133

## 134 **2 Methods**

135

### 136 2.1 Study area

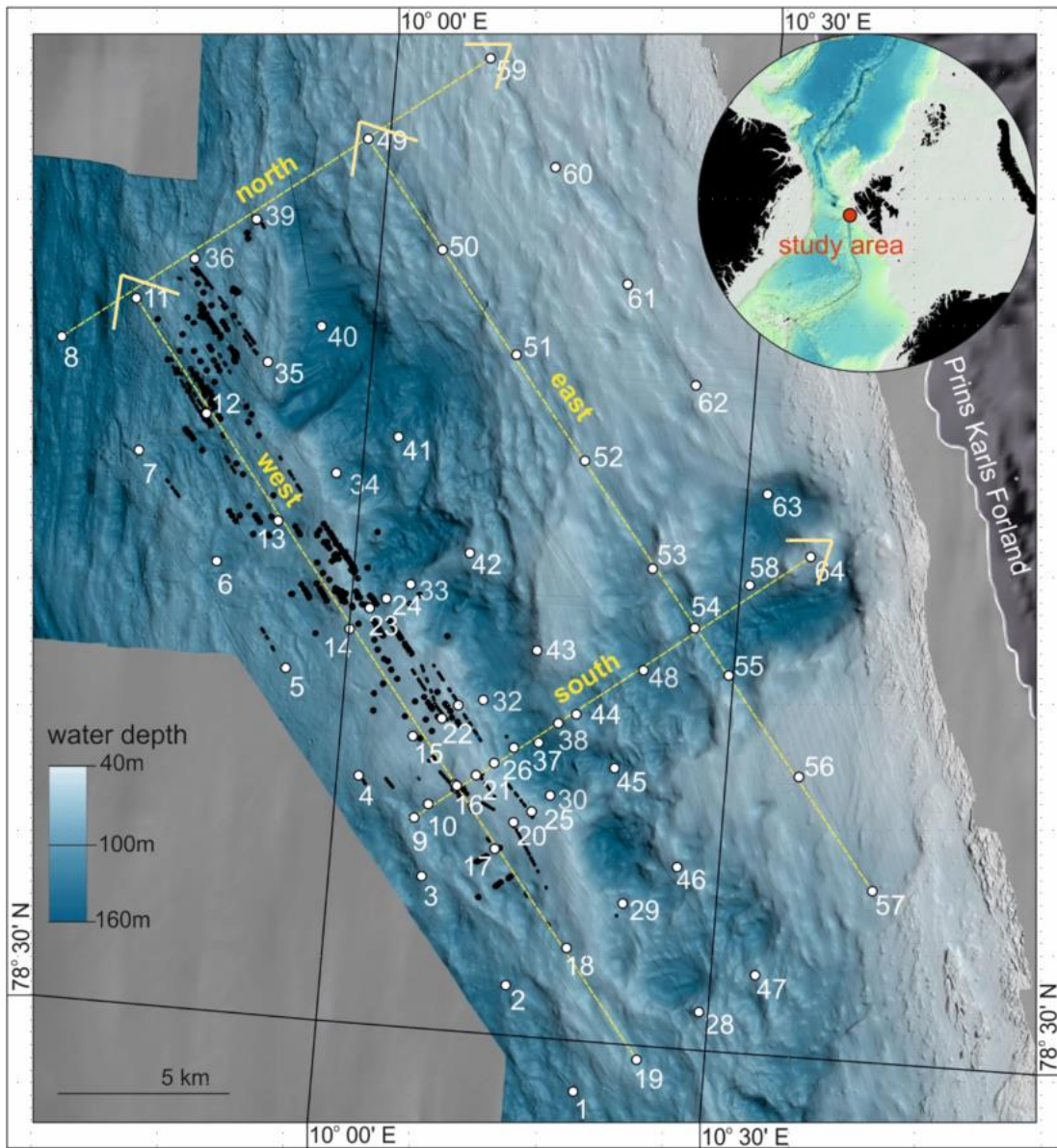
137 Our study area (423 km<sup>2</sup>; 50 - 120 m water depth) is located west of PKF (Fig. 1). The  
138 seafloor in this area is complex and characterized by abundant depressions and a sequence of  
139 pronounced end moraine ridges: the Forlandet moraine complex (Landvik et al., 2005). Several  
140 hundred methane flares were found during the present study and previous expeditions (e.g.  
141 Sahling et al., 2014 and references therein). Similar to the adjacent shelf break, gas seepage is not  
142 related to pockmarks or other fluid leakage related structures and the origin of the methane  
143 remains unconfirmed (Westbrook et al., 2009; Berndt et al., 2014). Although hydrates have never  
144 been recovered in the area and seismic evidence of gas hydrates is missing, sediment cores drilled  
145 outside PKF contained freshwater presumably originating from dissociated gas hydrates  
146 (Wallmann et al., 2018). Previous studies also suggest that free gas may originate from gas  
147 hydrate dissociation deeper on the continental slope (>300 m) where gas hydrates have been  
148 found (Sarkar et al., 2012) and migrate along the permeable zones towards the shelf (Westbrook  
149 et al., 2009). An alternate hypothesis is that glacial rebound at the beginning of the Holocene  
150 resulted in gas hydrate dissociation, which allowed for the formation of shallow gas pockets that  
151 continue to release methane into the water column (Portnov et al., 2016; Wallmann et al., 2018).

152 The water masses and circulation in the study area are controlled to a large extent by the  
153 interaction of coastal processes on the shelf with the West-Spitsbergen Current (WSC) that  
154 circulates northward along the shelf break as the northernmost extension of the North-Atlantic  
155 Current, transporting AW into the Arctic Ocean. The core of the WSC is at 250-800 meters water  
156 depth (Perkin and Lewis, 1984) and the stream follows the slope of the continental margin  
157 (Aagaard et al., 1987). By bringing large amounts of salt and heat, it affects the water column  
158 structure in the entire area. Other currents in the area are the East Spitsbergen Current (ESC) that  
159 advects Arctic waters into the region, and the coastal surface current, associated with the West  
160 Spitsbergen Polar Front (Nilsen et al., 2016). Local scale physical processes affecting water mass  
161 circulation include exchange of water masses between the WSC and shelf waters due to  
162 instability of the WSC core and resulting eddies (Teigen et al., 2010; Hattermann et al., 2016;  
163 Appen et al., 2016); as well as wind forcing and resulting upwelling events (Berge et al., 2005;  
164 Cottier et al., 2007).

165

## 166 2.2 Survey design

167 We conducted research expeditions with the R/V Helmer Hansen in the study area during  
168 three consecutive years on dates when conditions at sea are the most favourable for collecting  
169 data: 25-27 June 2014 (hereafter, June-14), 01 – 03 July 2015 (July-15), 02 – 04 May 2016 (May-  
170 16). Each year we visited 64 hydrographic stations. Stations were positioned in a grid, evenly  
171 distributed 2-3 km apart, closer above presumed dense clusters of methane flares based on  
172 ecosounder observations (Fig. 1). The total area covered by the grid was approximately 14x26  
173 km. Stations were sampled consecutively without any break (i.e., the entire grid was sampled  
174 within a time frame of ~72 hours). We collected hydrocast data from each station including  
175 continuous measurements of conductivity, temperature, depth (CTD), and sampled the water  
176 column at discrete depths for subsequent dissolved methane concentration measurements (see  
177 details in section 2.4). The entire grid was subsampled within 3 days during each survey.  
178 Underway hydro-acoustic scanning of the water column was performed to acquire information on  
179 gas flares (section 2.3). Ship-mounted meteorological instruments continuously recorded air  
180 temperature, atmospheric pressure, wind speed and direction. Furthermore, atmospheric methane  
181 mixing ratios were recorded continuously with a Cavity Ring-Down Spectrometer (CRDS,  
182 PICARRO G2401) with an air intake at 22.4 m above sea level (Platt et al., 2018).



184

185

186 **Figure 1.** Bathymetric map of the study area with 64 hydrographic stations (white dots) for  
 187 oceanographic measurements west of the Svalbard archipelago (overview map). Black dots  
 188 indicate locations of methane seeps detected on echograms during all three surveys. Yellow  
 189 dashed arrows indicate transects shown in Fig. 4. Bathymetry data were acquired on board with a  
 190 Kongsberg Simrad EM 300 multibeam echo sounder (frequency of 30 kHz).

191

192 2.3 Hydroacoustic data acquisition and gas flux calculations



193  
194 Gas bubbles in the water column were detected as acoustic signatures (flares) with a  
195 Kongsberg Simrad EK60 single beam echosounder system. This system is primarily designed for  
196 the fishery industry, but is also used to detect gas bubbles in the water column (Ostrovsky et al.,  
197 2008; Nikolovska et al., 2008). Data were acquired at 38 kHz as this is the most appropriate  
198 frequency to detect gas bubbles of sizes expected for cold seeps (Greinert et al., 2006). We used  
199 the FlareHunter program (Velooso et al., 2015) to distinguish flares from other echo signals such  
200 as fish, seafloor, and interference artifacts, and calculated flow rates from echosounder  
201 backscatter based on beam compensated Target Strength (TS, dB) in a 5-10-meter layer above the  
202 seafloor. We report free gas flow rates as mean values calculated from seven different bubble  
203 rising speed models (BRSMs). The relative uncertainty between BRSM estimates is 16 %  
204 (Velooso et al., 2015). The total length of our survey line was 408 km in June-14, 427 km in July-  
205 15, and 300 km in May-16. Accounting for water depth and the resulting beam width radii of 5 –  
206 20 m, the area of the seafloor investigated by the echosounder was 5.5 km<sup>2</sup> in June-14 and July-  
207 15, and 3.8 km<sup>2</sup> in May-16, which amounts to ~1 % of the total study area. Since the fraction of  
208 the study area covered by echosounder was small and slightly different between the three  
209 surveys, we applied Kriging interpolation to scale up estimates over the entire study area and thus  
210 facilitate comparison (details are provided in Appendix A and Fig. A1). Note that Fig. 2 shows  
211 observed flow rates of single sources. For comparison with other studies we present flow rates for  
212 the whole study area in Table 1 calculated as: (i) integrated over the entire area volumetric flow  
213 rate (L min<sup>-1</sup>); (ii) converted into mass flow rate (t y<sup>-1</sup>) using the ideal gas law and accounting for  
214 the average depth within each cell; and (iii) mean flux averaged over the whole area (mmol m<sup>-2</sup> d<sup>-1</sup>),  
215 converted from mass flow rate using the molecular weight of methane and divided by the  
216 survey area (423 km<sup>2</sup>).

#### 217 2.4 CTD profiling and water sample analyses

218 Vertical profiles of seawater temperature, salinity and pressure were recorded with a SBE  
219 911 plus CTD probe at a rate of 24 Hz. The probe was mounted on a rosette including 12 5-litre  
220 Niskin bottles. The Niskin bottles were closed during the up-cast (at speed of 1 m s<sup>-1</sup>). For  
221 analysis of hydrographic profiles, only down-casts were considered. Water samples were taken at  
222 5, 15 and 25 meters above the seafloor and below the sea surface, and an additional two samples

223 were collected at evenly spaced depth levels between 25 m above the seafloor and 25 m below  
224 the sea surface. In total, eight depths were sampled during all surveys.

225 Immediately upon recovery, sub-samples from the Niskin bottles were collected through  
226 silicon tubing into 60 ml plastic syringes (June-14) or 120 ml serum glass bottles (July-15, May-  
227 16) with rinsing by 2 – 3 overflow volumes. Syringes were closed with a 2-way valve and serum  
228 bottles were crimp-sealed with butyl rubber septa. 5 ml N<sub>2</sub> headspace was added to the syringes  
229 and serum bottles. Syringes/serum bottles with headspace were vigorously shaken for two  
230 minutes to allow the headspace N<sub>2</sub> to equilibrate with the dissolved methane in the water sample.  
231 Headspace methane mixing ratios were determined by gas chromatography (GC). During the  
232 June-14 survey a ThermoScientific FOCUS GC equipped with a flame ionization detector (FID),  
233 and a Resteck 2 m packed column HS-Q 80/100 with hydrogen (40 ml min) as a carrier gas was  
234 used. During the July-15 and May-16 surveys a ThermoScientific Trace 1310 GC equipped with  
235 an FID, and a Restek 30 m Alumina BOND/Na<sub>2</sub>SO<sub>4</sub> column with hydrogen as a carrier gas (40  
236 ml min<sup>-1</sup>) was used. The column temperature was held constant at 40°C. The systems were  
237 calibrated with external standards (2 ppm and 30 ppm in June-14 (Air Liquide); 10 ppm, 50 ppm,  
238 and 100 ppm in July-15 and May-16 (Carbagas). Finally, water column methane concentrations  
239 were calculated from headspace methane mixing ratios according to Wiesenburg & Guinasso  
240 (1979) with consideration of salinity, sample temperature and ambient atmospheric pressure.

## 241 2.5 Calculations of water column methane content

242 To account for the uneven bathymetry (bottom depths of 50 to 120 m), when comparing  
243 bottom, intermediate and surface waters, we divided the water column in three layers (Fig. B1):  
244 (1) a bottom layer (0-15 meters above seafloor), (2) an intermediate layer (15 meters above  
245 seafloor to 20 m water depth; the upper boundary roughly follows the depth of the pycnocline  
246 during the July-15 survey, which we determined as a function of the Brunt–Väisälä frequency,

247 see Fig. 5) and (3) a surface layer (20 m water depth to sea surface). Detailed calculations of the  
248 methane content (in mol) within the study area can be found in Appendix B.

## 249 2.6 Calculations of the sea-air methane flux

250 The sea-air methane flux was calculated according to Wanninkhof et al. (2009), using the  
251 methane concentration ( $C_o$ ) at the ocean surface in presumed equilibrium with the atmosphere,  
252 and measured methane concentration in the well-mixed surface layer. The gas transfer was  
253 calculated using the wind speed at 10 m above the sea surface, derived from the wind speed  
254 measured by the ships' anemometer at height 22.4 m after Hsu et al. (1994). The Schmidt number  
255 is the non-dimensional ratio of gas diffusivity and water kinematic viscosity, and was defined as  
256 677 in accordance with Wanninkhof et al. (2009). The  $C_o$  was calculated with the Bunsen  
257 solubility of methane in seawater according to Wiesenburg and Guinasso (1979). Detailed  
258 description of flux calculations with used equations can be found in Appendix C.

## 259 2.7 Modelling of water mass properties and particle release experiments

260 To study seasonal variations in water mass properties and circulation and to scale up our  
261 observations to a full year, we used a high-resolution regional ocean sea ice model. A more  
262 detailed description and validation of the Svalbard 800 m horizontal resolution model (the S800-  
263 model hereafter) can be found in Albretsen et al. (2017), Hattermann et al., (2016) and Crews et  
264 al. (2018). Briefly, the S800-model provides hindcast ocean sea ice simulations for the Svalbard  
265 and the Fram Strait region (see also Fig. 1 in Hattermann et al. (2016) for overview of the spatial  
266 extent of the model domain) based on the Regional Ocean Modelling System (ROMS,  
267 Shchepetkin and McWilliams, 2005) and a coupled sea ice component (Budgell, 2005). Boundary  
268 conditions are provided by a 4 km pan-Arctic setup (A4-model). The bathymetry is based on the  
269 ETOPO1 topography (Amante, 2009). Vertically, the model is discretized into 35 levels with a  
270 layer thickness of less than 1 m near the surface over the continental shelf. The S800-model is  
271 initialized and forced with daily averages from the A4-model, for which boundary conditions and  
272 forcing fields are based on reanalyses (Storkey et al., 2010). Atmospheric forcing is provided by  
273 ERA-interim reanalysis (Dee et al., 2011) and climatological river input from major rivers in the  
274 area, including freshwater runoff from the Svalbard archipelago (details in Hattermann et al.,

275 2016). The S800-model was initialized from January 2005, and the data shown in this study are  
276 based on model runs from July 2005 to July 2010.

277 From the daily 3D model output, we computed monthly climatologies (averaging data  
278 from all January months, all February months, etc.) of potential temperature and salinity as well  
279 as horizontal mean and eddy kinetic energy. The latter being defined as the average of the daily  
280 deviations from the climatological mean (see also Hattermann et al. (2016) for detailed  
281 definitions). These results were then horizontally averaged on a subset of  $41 \times 56$  grid points that  
282 corresponded to the geographic area of the survey between CTD stations 1, 8, 57, 59, 64 (Fig.1;  
283 also red polygons in Fig. 10).

284 To investigate seasonal features of methane dispersion and displacement in the study area,  
285 we conducted numerical experiments by simulative release of neutrally buoyant Lagrangian  
286 drifters (hereafter termed particles) that were advected by the 3D model velocity field. Those  
287 trajectories were computed using Lagrangian particle tracking algorithm TRACMASS (Döös et  
288 al. 2017) based on the daily S800-model output (see Hattermann et al. (2016) for details). We  
289 released particles from the polygon where the most intense seeps were observed during the  
290 surveys. The polygon enclosed CTD stations 3 (113 m water depth), 4 (103 m water depth), 15  
291 (91 m water depth), and 17 (97 m water depth) (Fig. 1). Due to varying water depths at these  
292 stations, we chose to release particles from uniform depths between 80 and 100 m. Particles were  
293 released every day throughout the 5-year long simulation period and were tracked until a  
294 maximum lifetime of ten days. Monthly maps of particle density distributions were computed  
295 from the particle positions five days after their release within a respective month. These  
296 histograms of particle distributions were computed by bin-counting particle positions on the  
297 S800-model lattice, normalized by the total number of particles. To assess the residence of  
298 particles in the study region, we computed the fraction of particles that remained within the study  
299 region as a function of particle lifetime and month of the particle release. To assess the role of  
300 different water masses for the particle dispersion, we evaluated the ambient density from the  
301 S800-model output at the particle position after five days and computed a two-dimensional  
302 histogram of particle counts as a function of distance from the source and potential density of  
303 ambient water.

### 304 **3 Results and discussion**

#### 305 3.1 Controls on flare abundance and methane flux from sediments

306  
307 We observed the densest flare cluster in the western and north-western part of the study  
308 area (Fig. 2). This cluster was venting free gas during all cruises. In contrast, there was a  
309 difference in flare density between surveys in the southern part of the study area, with the highest  
310 flare density during the June-14 survey, and much lower densities during the July-15 and May-16  
311 surveys. In total, we counted 225 individual flares in June-14, 208 in July-15 and only 92 during  
312 the May-16 survey. The estimated gas flux from individual flares ranged between 20 and 600 ml  
313 min<sup>-1</sup> (Fig. 2). As a consequence of the decreasing flare density from June-14 to May-16, the  
314 calculated total volumetric gas flow rate over the surveyed area was larger for June-14 (900 L  
315 min<sup>-1</sup>) than for the July-15 (665 L min<sup>-1</sup>) and May-16 surveys (540 L min<sup>-1</sup>) (Table 1).

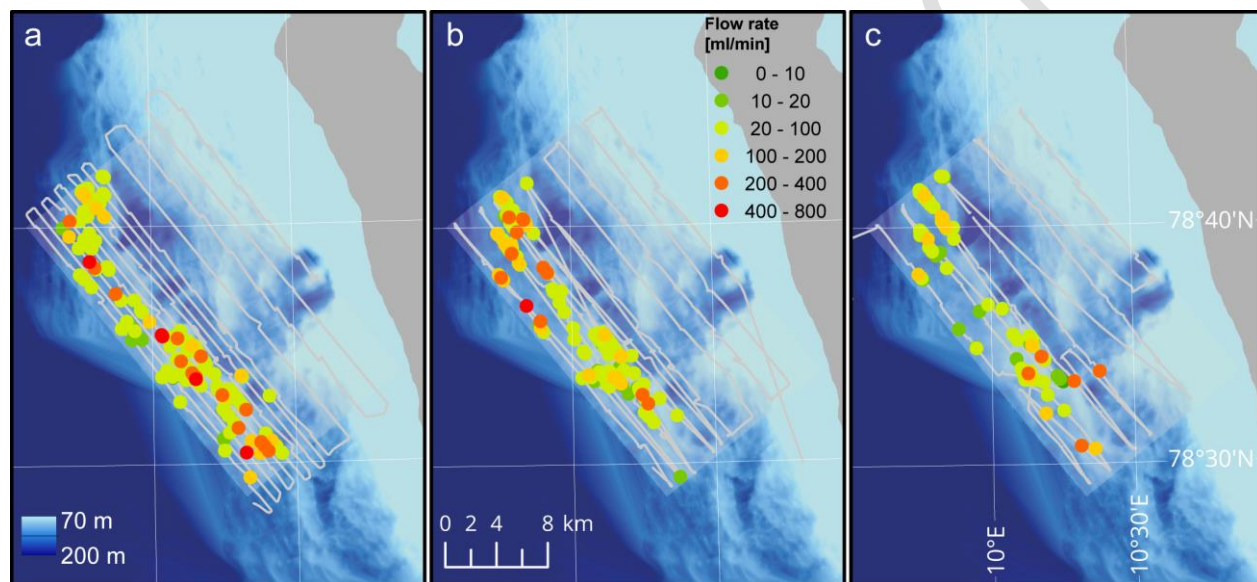
316 We carefully checked for factors that may have potentially biased our estimates. The  
317 May-16 survey was substantially shorter in distance (~70% compared to June-14 and July-15),  
318 decreasing the confidence in scaling up our observations to the entire area. Yet, the western part  
319 of the study area, where we always observed the highest flare density, was investigated during all  
320 three surveys. Considering only this area, we could still identify a substantial decrease in both  
321 flare density and volume flux. Consequently, artefacts from the scaling up the observations made  
322 during surveys of different distance cannot explain the observed differences in seepage activity.

323 Temporal variability in the activity of seafloor methane seeps has been reported  
324 previously (e.g. Greinert et al., 2006; Klauke et al., 2010; Kannberg et al., 2013). Römer et al.  
325 (2016) investigated a cold seep offshore Vancouver Island at 1250 m water depth and suggested  
326 that the pressure change of 1.9 dbar between low and high tide affected seepage activity with  
327 increasing gas flux during falling tides. However, our survey period lasted for ~3 days, i.e. ~6  
328 tidal cycles, so that potential forcing by tides should be equalised and tides cannot be the reason  
329 for differences in seepage activity between the surveys.

330 Variability in gas flux in our study area (highest in June-14, lower in July-15 and lowest  
331 in May-16) follows observed between-survey differences in bottom water temperature (Fig. D1).  
332 This was highest in June-14 ( $3.63 \pm 0.2^\circ\text{C}$ ), lower in July-15 ( $3.49 \pm 0.2^\circ\text{C}$ ) and the lowest in May-

333 16 ( $1.77 \pm 0.1$  °C). Indeed, it has been proposed that seasonal fluctuations in bottom water  
 334 temperature modulate seepage activity off Svalbard, but from gas hydrate bearing sediments at  
 335 the termination of the gas hydrate stability zone (Berndt et al., 2014). However, gas hydrates have  
 336 never been found in our study area, which is at ~200 m shallower water depth than that of gas  
 337 hydrate stability limit (>300 m water depth), so that we can only speculate about the mechanisms  
 338 of a potential temperature control on seepage activity. Nevertheless, potentially modulating  
 339 effects of bottom water temperature would imply seasonal fluctuations in seepage activity in our  
 340 study area.

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342

343

344 **Figure 2.** Flow rates from single sources (flares) during June-14 (a), July-15 (b) and May-16 (c)  
 345 surveys. Coloured circles indicate gas flow rates in  $\text{ml min}^{-1}$  from individual flares on the  
 346 seafloor. The grey line represents the ship track and echosounder beam coverage.

347

348 **Table 1.** Methane fluxes from sediments in different surveys

Survey	Total volumetric flow rate in the area ( $\text{L min}^{-1}$ )	Total mass flow rate in the area ( $\text{t y}^{-1}$ )	Average methane flux from sediments ( $\text{mmol m}^{-2} \text{d}^{-1}$ )

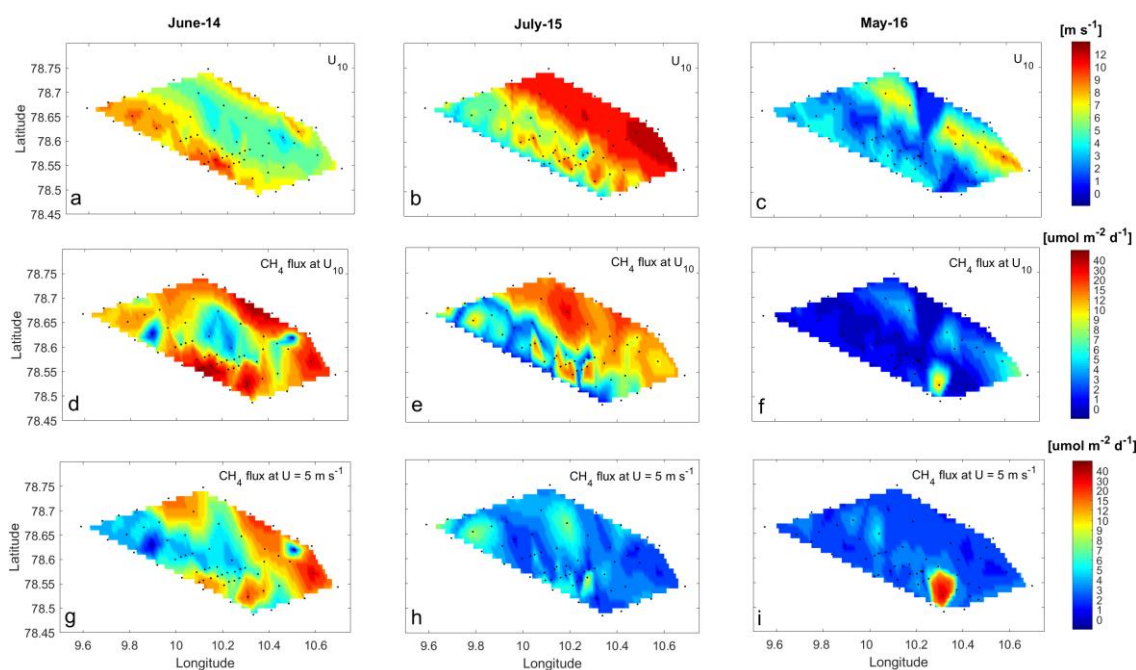
<b>June-14</b>	899	3774	1.53
<b>July-15</b>	665	3004	1.21
<b>May-16</b>	542	2356	0.96

### 349 3.2 Controls of sea-air methane flux

350 The highest sea-air methane flux of  $15 \mu\text{mol m}^2 \text{d}^{-1}$  was observed during the June-14  
351 survey, a lower flux of  $11 \mu\text{mol m}^2 \text{d}^{-1}$  was observed in July-15, and the lowest flux of only 2  
352  $\mu\text{mol m}^2 \text{d}^{-1}$  was observed during the May-16 survey (Fig. 3 d-f). The temporal pattern of  
353 atmospheric methane mixing ratios was the opposite of that of the flux, i.e. we found the lowest  
354 mixing ratios in June-14 ( $1902 \pm 0.52$  ppb), higher during the July-15 ( $1917 \pm 3.30$  ppb) and the  
355 highest during the May-16 survey ( $1955 \pm 25.4$  ppb) (data given as average  $\pm$  standard deviation  
356 of all observations during each survey). Thus, the atmospheric mixing ratio of methane was one  
357 of the main controls on sea-air fluxes resulting in a suppressed flux in case of higher atmospheric  
358 methane values (e.g. lower fluxes in May-16 compared to the highest encountered atmospheric  
359 methane mixing ratios). A further key control on sea-air methane fluxes is the concentration of  
360 methane in the well-mixed surface waters, which was  $9 \text{ nmol L}^{-1}$  in June-14, and  $3 \text{ nmol L}^{-1}$   
361 during the July-15 and May-16 surveys (Table 2). Despite the similar surface water  
362 concentrations in July-15 and May-16, sea-air methane fluxes were 5 times higher in July-15 than  
363 in May-16. This can be explained by the wind speed, which was comparably low and varied very  
364 little during the June-14 ( $4\text{-}8 \text{ m s}^{-1}$ ) and May-16 surveys ( $1\text{-}6 \text{ m s}^{-1}$ ), but increased from calm  $4\text{-}6$   
365  $\text{m s}^{-1}$  to strong  $10\text{-}12 \text{ m s}^{-1}$  towards the end of the 3-day July-15 survey (Fig. 3b). Generally, the  
366 differences between the atmospheric methane mixing ratio and surface water methane content as  
367 well as wind speed determine the variation in average sea-air flux. However, we argue that wind  
368 speed plays the most important role in our study area with respect to sea-air methane fluxes. High  
369 wind speeds can intensify efflux to the atmosphere even if the surface water methane  
370 concentration is relatively low as long as the surface waters are supersaturated with respect to the  
371 atmosphere.

372 To further test how the wind speed affects sea-air methane flux, we determined what the  
373 flux would have been if the wind speed had been a  $5 \text{ m s}^{-1}$  throughout all surveys. In other words,  
374 we used the observed values of surface water methane concentrations and atmospheric methane  
375 mixing ratios measured during each survey, but instead of the measured wind data, we calculated

376 fluxes for a constant wind speed of  $5 \text{ m s}^{-1}$ , which is the climatological average wind speed for  
 377 late spring to early summer in our study area (The Norwegian Meteorological Institute,  
 378 [www.yr.no](http://www.yr.no)). The meteorological mean was lower than the measured wind speed in June-14 and  
 379 July-15, but higher than the measured wind speed in May-16. Therefore, our flux calculations  
 380 with the mean values produced lower flux values for the June-14 ( $10 \mu\text{mol m}^2 \text{ d}^{-1}$ ) and July-15 ( $4$   
 381  $\mu\text{mol m}^2 \text{ d}^{-1}$ ) surveys, but higher values for the May-16 ( $3.5 \mu\text{mol m}^2 \text{ d}^{-1}$ ) survey (Fig. 3 g-i). This  
 382 comparison between sea-air methane flux with actual measured and constant wind speeds  
 383 highlights the importance of wind speed in modifying methane emission to the atmosphere in our  
 384 study area.



385  
 386  
 387 **Figure 3.** Wind speed measured at 22.4 m above sea level (upper panel: **a, b, c**), Methane flux at  
 388 the air-sea interface at measured wind speed  $u_{10}$  (mid panel: **d, e, f**) and methane flux at the air-



389 sea interface at constant wind speed  $U=5 \text{ m s}^{-1}$  for May-July (lower panel: **g, h, i**), for the entire  
390 grid and the three surveys.

### 391 3.3 Controls of water column methane content

392  
393 The water column above active methane flares in the study area was divided into three  
394 layers in order to estimate differences between methane content in the bottom 15 m, where  
395 presumably most of released methane dissolves; the surface 20 m which roughly corresponds to  
396 the thickness of the well mixed surface layer in summer and from which outgassing most of  
397 methane to the atmosphere occurs; and the intermediate layer between the bottom and surface  
398 layers, which is the thickest and presumably accumulates most of the released methane. When  
399 comparing different layers, the highest methane concentrations were found in bottom layer as  
400 expected. However, in all surveys the overall highest methane content was found in the  
401 intermediate layer because it contains the highest volume of water (extends through the largest  
402 depth interval). When comparing different surveys, we observed the highest total methane

403 content in June-14 ( $23 \times 10^5$  mol), lower in July-15 ( $15 \times 10^5$  mol) and lowest during the May-16  
404 ( $14 \times 10^5$  mol) survey (Table 2).

405

406 **Table 2.** Average dissolved methane concentrations and content in different layers during each  
407 survey

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<b>Layer</b> →	<b>Surface</b>	<b>Intermediate</b>	<b>Bottom</b>	<b>Total</b>
<b>Survey</b> ↓	<b>(surface-20 m water depth)</b>	<b>(variable depth depending on water depth)</b>	<b>(bottom-15 m above the bottom)</b>	
	Average methane concentrations (nmol L <sup>-1</sup> )			
<b>June-14</b>	9.4	55.4	92.3	
<b>July-15</b>	3.1	31.9	70	
<b>May-16</b>	3.2	26.6	61.3	
	Average content (×10 <sup>-3</sup> mol m <sup>-2</sup> )			
<b>June-14</b>	0.17	3.79	1.39	5.35
<b>July-15</b>	0.06	2.36	1.04	3.46
<b>May-16</b>	0.07	2.32	0.91	3.30
	Total content in the surveyed area (×10 <sup>5</sup> mol)			
<b>June-14</b>	0.73	16	5.87	23
<b>July-15</b>	0.26	10	4.40	15
<b>May-16</b>	0.28	9.8	3.85	14
	Total mass of methane in the surveyed area (t)			
<b>June-14</b>	1.17	25.73	9.41	36.31
<b>July-15</b>	0.43	16.00	7.05	23.50
<b>May-16</b>	0.44	15.77	6.17	22.38

408           The change in dissolved methane content in the water column between the surveys is  
409 similar to the trend in the number of observed flares and the volume of released gas, and, to a  
410 smaller extent, the sea-air methane flux. Although the correlation between the amount of released

411 methane and its content in the water is anticipated, there are number of processes that we did not  
412 measure, some of which could alter the methane content in the entire water column, and some, in  
413 surface waters alone.

414 One of these processes is aerobic methane oxidation (MOx), which leads to methane  
415 undersaturation of deep waters in the entire ocean (Reeburgh, 2007). During MOx, methane is  
416 removed from the water column when it is consumed by bacteria, which use methane as a source  
417 of carbon and energy. To test how important the role of MOx is in the removal of methane from  
418 the system, we used MOx rates reported for the regions near our study area. Gentz et al. (2014)  
419 reported MOx rate of  $0.8 \text{ nmol L}^{-1} \text{ d}^{-1}$  in bottom waters and  $0.2$  in surface waters in the water  
420 column above methane flares with absolute depth of  $\sim 250$  m, while Steinle et al. (2015) found  
421 higher rates of  $2 \text{ nmol L}^{-1} \text{ d}^{-1}$  in bottom water alongside lower rates of only  $0.1 \text{ nmol L}^{-1} \text{ d}^{-1}$  in  
422 surface waters above methane flares with an absolute water depth of 360 m. After vertical and  
423 horizontal integration of these estimates over our area, we found that less than 10% of the  
424 released methane in our study area per day is likely to be removed from the system through MOx,  
425 suggesting that this process does not play a major role in the removal of methane injected from  
426 sediments at this site.

427 In addition to methane oxidizing bacteria, benthic chemosynthetic autotrophs consume  
428 methane supplied from sediments. Biological systems in the vicinity of cold seeps are high in  
429 biomass and number of species, thereby resembling some aspects of hydrothermal vents (e.g.  
430 Boetius and Suess, 2004), although cold seep systems are not driven by hydrothermal or volcanic  
431 activity. Both types of habitats are enriched in reduced compounds such as methane and sulphide,  
432 though methane is typically more important at cold seeps, where sulphide is generated through  
433 anaerobic methane oxidation. At PKF, methane associated tubeworms and bacterial mats were  
434 observed but these appeared to be far less dense than at other highly active seeps such as Håkon  
435 Mosby Mud Volcano or Hydrate Ridge, possibly because of predation pressure (Carney, 1994;  
436 Sahling et al., 2003).

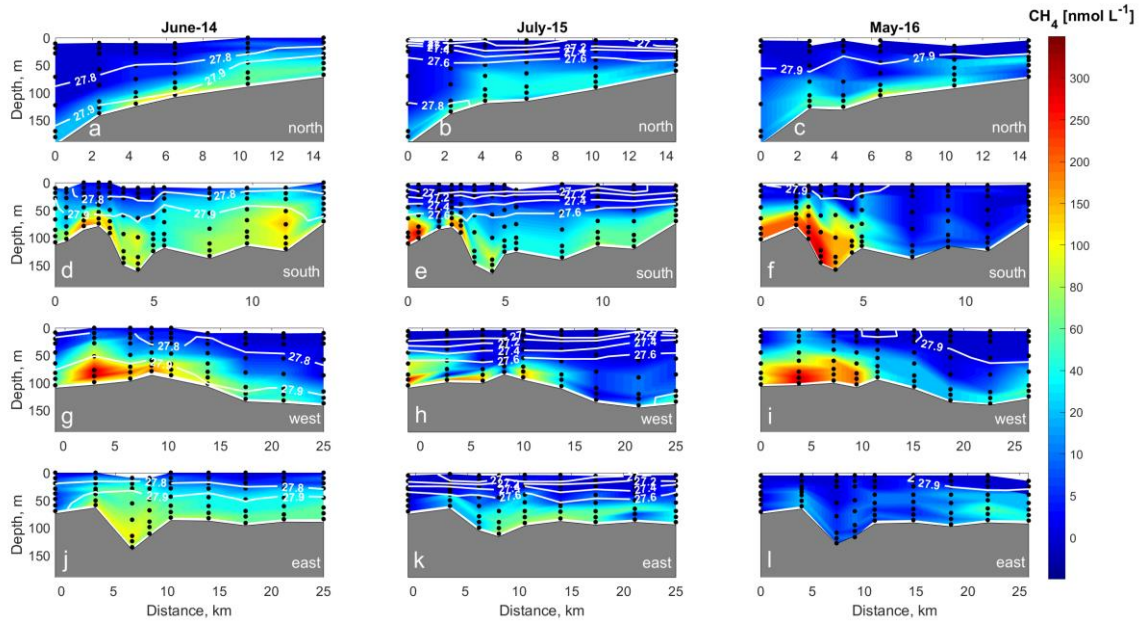
437 Another process mediating methane content in the water column is aerobic methane  
438 production by microbes under phosphorus limiting conditions (Karl et al., 2008). In the oceanic  
439 interior, this process leads to methane supersaturation in the surface water column above the  
440 pycnocline (Reeburgh, 2007). Such methane supersaturation in surface waters was found in the

441 Fram Strait to the west from our study area, but only reached maximum concentrations of 9 nM  
442 at 10-20 m depth (Damm et al., 2015). We observed only one case of isolated high surface  
443 methane concentration (of 20 nmol L<sup>-1</sup>) during the June-14 and May-16 surveys, but in most  
444 cases surface concentrations were close to atmospheric equilibrium, thus we assume that in our  
445 study area the methane contribution from this process is of low importance.

446 These three biological processes are important on the scale of entire ocean but are minor  
447 mediators of methane content in our study area, which experiences rapid methane injection into  
448 the system at the seabed and methane concentrations hundreds of times higher than the average  
449 oceanic concentrations. For example, considering a total injection of methane from sediments of  
450  $5.2 \times 10^5$  mol d<sup>-1</sup> (averaged over the three surveys), a loss through sea-air gas exchange of  $0.04 \times$   
451  $10^5$  mol d<sup>-1</sup>, and a MOx rate of  $0.58 \times 10^5$  mol d<sup>-1</sup> (based on estimates from Gentz et al. (2014)  
452 and Steinle et al. (2015) for nearby waters), the resulting amount of methane in the water column  
453 would be  $4.6 \text{ mol} \times 10^5$  mol d<sup>-1</sup>. Our total methane content averaged over the three surveys is  $17 \times$   
454  $10^5$  mol, which is 3.8 times higher than the resulting content, implying a residence time of  
455 methane in the study area of about 3.8 days. However, methane is likely transported beyond our  
456 survey area during this time through transport by lateral water movement (section 3.5). To see  
457 how efficient this transport is and what affects it in our study area, we look further into vertical  
458 and horizontal distribution of methane in different surveys.

### 459 3.4 Controls of the vertical distribution of dissolved methane

460  
461 Highest dissolved methane concentrations were found in the bottom layer (> 300 nmol L<sup>-1</sup>)  
462 in the south-western part of the sampling area during all three surveys (Figs. 4d-i). Waters  
463 supersaturated with methane were found around flares from the seafloor up to 50 (July-15) and  
464 20 meters water depth (June-14, May-16). Methane supersaturated waters have methane  
465 concentrations > 3.7 nmol L<sup>-1</sup>, which would be in equilibrium with the atmosphere for a salinity  
466 of 35 at 0° C and atmospheric mole fraction of methane 1.9 ppb (average value for all three  
467 surveys) (Wiesenburg and Guinasso, 1979). In all three surveys, the intermediate layer methane  
468 concentration averaged over the entire area was only half of the bottom layer concentration, while  
469 the surface water concentrations were 25 times lower than the bottom layer concentrations.



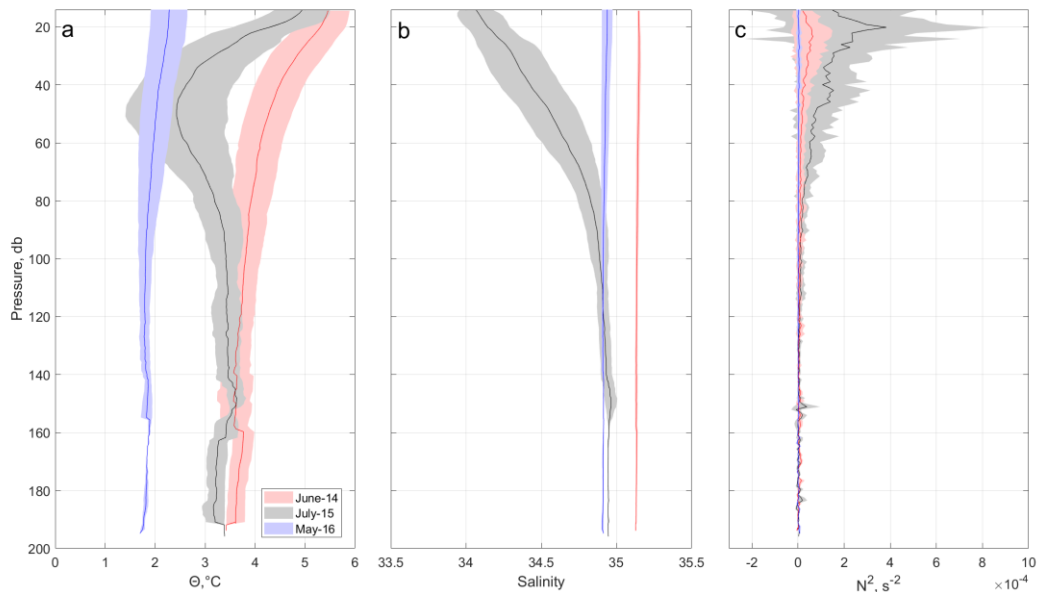
470  
 471 **Figure 4.** Distribution of dissolved methane along four transects (north (**a, c, d**), south (**d, e, f**),  
 472 west (**g, h, i**) and east (**j, k, l**); c.f. Fig. 1 for location and direction of each transect). Seawater  
 473 density (in  $\text{kg m}^{-3}$ ) is indicated by white contour lines. Locations of discrete samples for methane  
 474 concentration measurements are indicated by black dots.

475 Our results show methane enriched bottom and intermediate waters, and surface water  
 476 which are only slightly supersaturated or close to atmospheric equilibrium. These results agree  
 477 well with earlier measurements near our study area (e.g. Gentz et al., 2014; Westbrook et al.,  
 478 2009; Mau et al., 2017), which showed high methane concentrations in bottom waters above  
 479 methane flares, and rapid decreases in methane concentrations towards the surface. This pattern  
 480 in vertical distribution can be explained by ongoing gas exchange between rising methane  
 481 bubbles and the surrounding seawater (e.g. McGinnis et al., 2006). This leads to continuous  
 482 replacement of methane in the bubbles with  $\text{N}_2$  and  $\text{O}_2$  from the seawater and methane  
 483 enrichment of seawater along the bubble ascent. Modelling approaches suggest that the bulk of  
 484 methane is already stripped out from rising bubbles close to the seafloor, so that bottom waters  
 485 become more enriched with dissolved methane (McGinnis et al., 2006). Bubbles observed close  
 486 to the surface are thus mostly comprised of  $\text{N}_2/\text{O}_2$ . Only bubbles of  $>20$  mm in diameter may still

487 contain 1% of their initial methane content at the surface, but such bubbles typically break apart  
488 during their ascent (McGinnis et al., 2006).

489 Vertical transport of dissolved methane that has already escaped bubbles has been  
490 proposed to be limited by water column vertical stratification, when a pycnocline acts as a barrier  
491 for vertical mixing of methane rich waters in strongly stratified waters (Gentz et al., 2014; Myhre  
492 et al., 2016). As a proxy for water column vertical stratification, we calculated the Brunt–Väisälä  
493 frequency ( $N^2$ ) in our study area (Millard et al., 1990), which generally peaked at ~20m water  
494 depth, and was the highest in July-15 ( $4 \times 10^{-4} \text{ s}^{-2}$ ), ~8 times lower during the June-14 survey  
495 ( $0.5 \times 10^{-4} \text{ s}^{-2}$ ) and near zero in the entire water column during the May-16 survey ( $0.1 \times 10^{-4} \text{ s}^{-2}$ )  
496 (Fig. 5c). In July-15 the observed strong stratification was formed by a temperature drop from 5.5  
497 °C at the surface to 3.5°C at 50 m water depth forming pronounced thermocline (Fig. 5a); and by  
498 a salinity increase from 34.1 at the surface to 34.9 at 100 m depth along a continuous halocline

499 (Fig. 5b). Conversely, in May-16 the water column was well-mixed, with almost uniform  
500 temperature and salinity with depth, and the near- absence of a pycnocline.



501

502 **Figure 5.** (a) Potential temperature ( $\Theta$ , °C), (b) salinity and (c) Brunt–Väisälä frequency  
503 ( $N^2$ ,  $s^{-2}$ ) averaged over all CTD stations for each survey with standard deviation shown as shaded  
504 error bars. Colours indicate: June-14 (red), July-15 (black) and May-16 (blue).

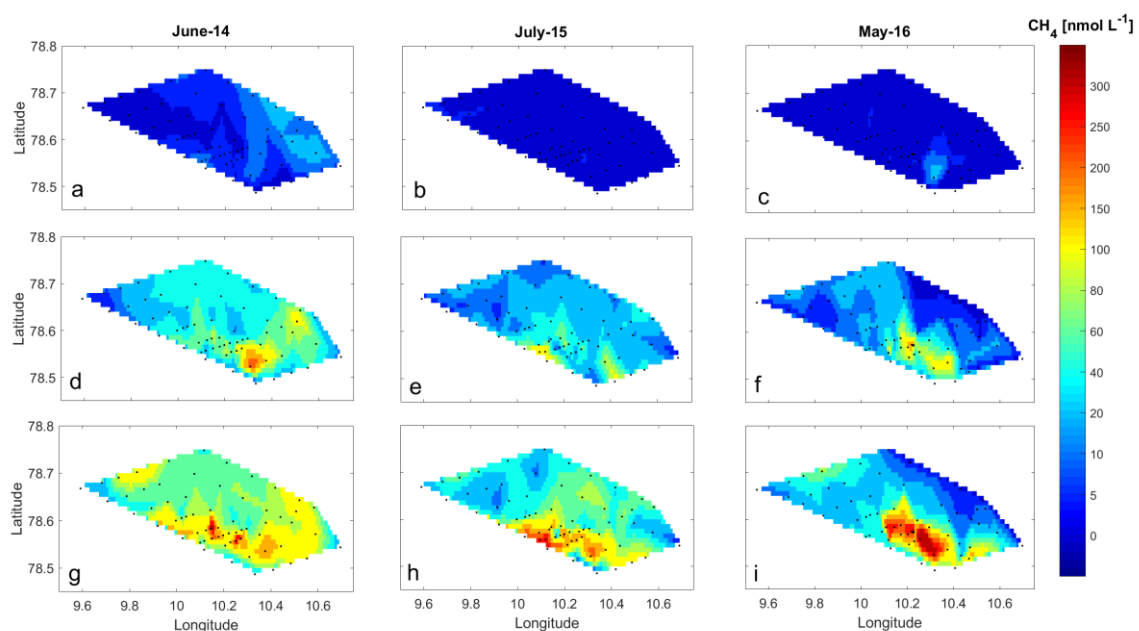
505 Despite the difference in stratification between the three surveys (Fig. 5c), the vertical  
506 distribution of dissolved methane (high bottom water methane concentrations and low surface  
507 water concentrations) was similar across all three surveys (Fig. 4). This indicates that methane  
508 released from the sediments and dissolved in seawater did not rise above 20-50 m water depth  
509 towards the sea surface, even in the absence of a pycnocline. Our findings thus suggest that water  
510 density stratification may not always play the principle role in the vertical distribution of  
511 dissolved methane in cold seeps areas, in contrast to the conclusions of previous studies in this  
512 area (Myhre et al., 2016; Gentz et al., 2014). Furthermore, our results do not show an influence of  
513 stratification on water column methane content or the sea-air gas flux.

514 3.5 Controls of horizontal distribution of dissolved methane



515 The horizontal distribution and patchiness of methane differed between the three surveys.  
516 During the June-14 survey we observed elevated dissolved methane concentrations in the bottom  
517 and mid-depth layers (Fig. 6d and g) spread over the entire survey area. In contrast, during May-  
518 16, methane concentrations were high (up to 400 nmol L<sup>-1</sup>) only above flares clustered in the  
519 south-western part of the area, and decreased considerably to < 40 nmol L<sup>-1</sup> within a few hundred  
520 meters away from the flares (Fig. 6i). Elevated methane concentrations also spread horizontally  
521 in July-15, but to a lesser extent than during the June-14 survey.

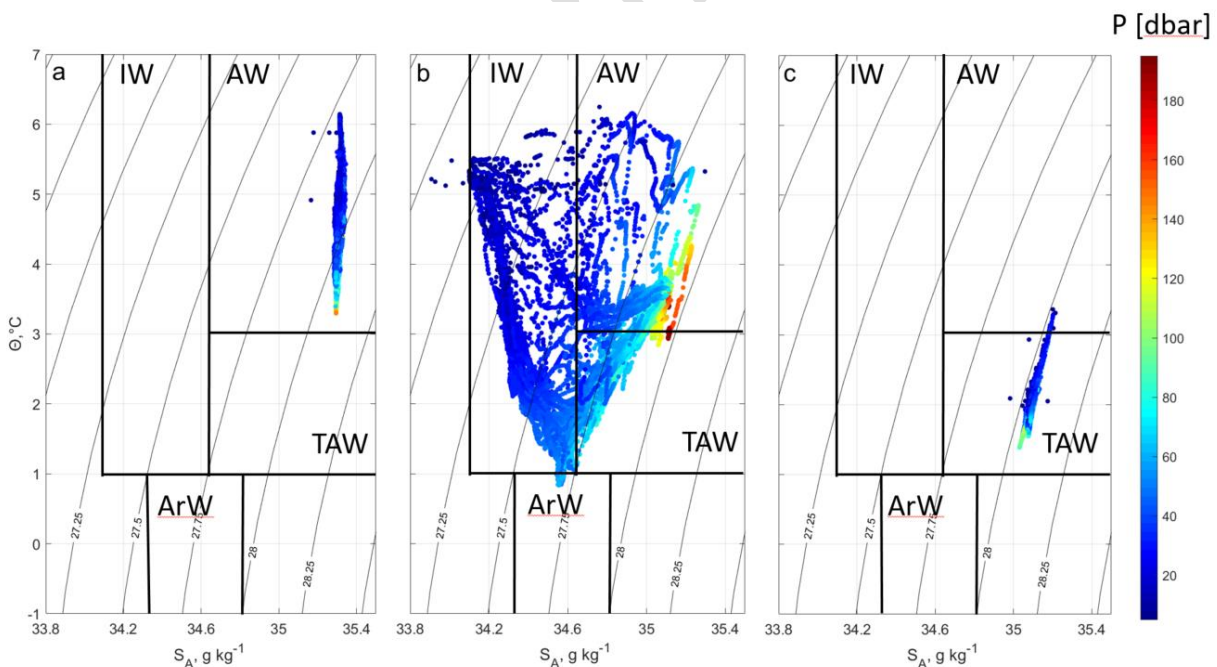
522 In the surface layer, methane concentrations were generally low and near the atmospheric  
523 equilibrium (Fig 6a-c). Some elevated surface methane concentrations (~20 nmol L<sup>-1</sup>) were  
524 observed at one station in the southeast part of the study area in June-14 and in the south-western  
525 part of the study area during the May-16 survey.



526  
527 **Figure 6.** Average methane concentrations in the surface (0-20 m, **a, b, c**), intermediate (20 m –  
528 15 m from the seafloor, **d, e, f**) and bottom water (within 15 m of the seafloor, **g, h, i**), layers for  
529 the entire grid during the three surveys as indicated above the figures.

530 High variability in water mass properties indicates that circulation during all surveys was  
531 controlled by several factors. We used the classification of water masses suggested by Cottier et  
532 al. (2005) for Svalbard fjords and adjacent shelf regions to describe the oceanographic setting in

533 our study area. During the June-14 survey we observed only warm and saline AW (temperature  $\Theta$   
 534  $>3^{\circ}\text{C}$ , absolute salinity  $S_A >34.65$ ) (Fig. 7a), brought to the study area with the WSC. In contrast,  
 535 water in July-15 was substantially colder and less saline (Fig. 7b), mainly comprised of AW, with  
 536 some Transformed Atlantic Water (TAW,  $1 < \Theta < 3^{\circ}\text{C}$ ,  $S_A >34.65$ ), and to the largest extent,  
 537 Intermediate Water (IW,  $\Theta >1^{\circ}\text{C}$ ,  $34 < S_A < 34.65$ ). IW originates from fjords and forms as AW  
 538 that cools over winter in fjords, and is freshened by glacial melt, sea ice melt and river runoff  
 539 during summer. IW can also be a mix of AW and Arctic Water masses (ArW,  $-1.5^{\circ} < \Theta < 1^{\circ}\text{C}$ ,  
 540  $34.3 < S_A < 34.8$ ) transported from the Northern Barents Sea around southern tip of Svalbard  
 541 with the ESC. During the May-16 survey (Fig. 7c), the water column mainly comprised TAW  
 542 with absolute salinity values similar to AW but with potential temperatures around  $1.5 - 3^{\circ}\text{C}$ ,  
 543 which is colder than the typical AW with temperature defined as above  $3^{\circ}\text{C}$ . There was a strong  
 544 presence of AW on the shelf and adjacent fjords in 2016 (F. Nilsen, pers. comm.). The core of the  
 545 AW in May is always above  $2.5^{\circ}\text{C}$  (Beszczynska-Möller et al., 2012). Our measured colder  
 546 seawater temperatures in the area could indicate that AW was cooled by the atmosphere or  
 547 surrounding waters, either locally or before it was advected from adjacent basins.

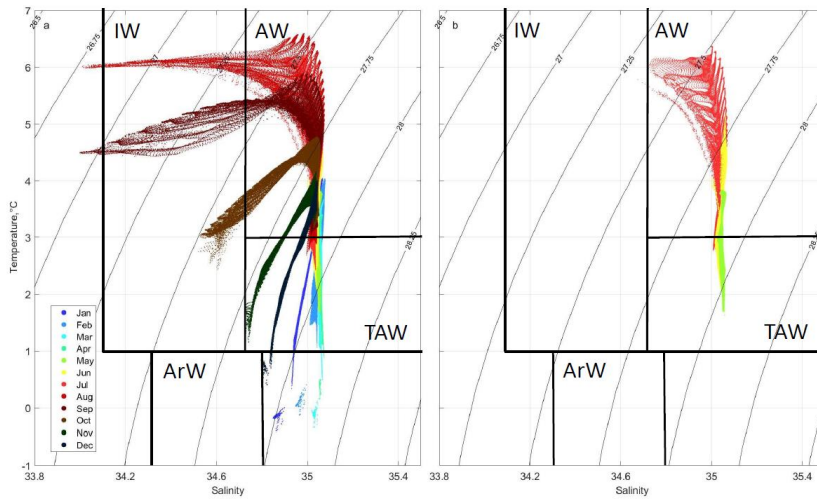


548  
 549 **Figure 7.** Potential temperature ( $\Theta$ ,  $^{\circ}\text{C}$ ) – absolute salinity ( $S_A$ ,  $\text{g kg}^{-1}$ ) diagrams for the June-14  
 550 (a), July-15 (b) and May-16 (c) surveys.  $\Theta$ ,  $^{\circ}\text{C}$  calculated according to the International  
 551 Thermodynamic Equation of Seawater (Fofonoff and Millard, 1983). Absolute salinity calculated

552 based on measured practical salinity, and is expressed in terms of g of salt per kg of water. Grey  
553 contours indicate isopycnals ( $\text{kg m}^{-3}$ ).

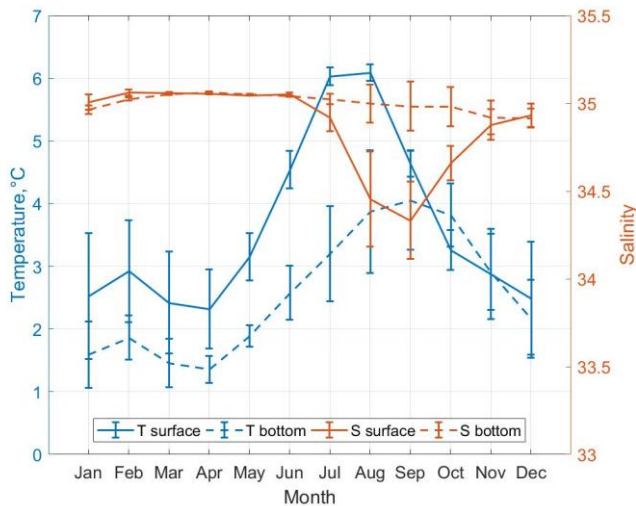
554 Seawater temperature and salinity modelled with the S800-model (Hattermann et al.,  
555 2016) for the study area indicate a shift from AW to IW properties towards summer and autumn  
556 months due to surface warming and freshening (Fig. 8), hence revealing the annual cycle of water  
557 mass formation. Salinity shows a seasonal cycle only at the surface, where it decreases from 35 in  
558 June to 34.4 in September (Fig. 9). Summer freshening of the surface results from freshwater  
559 runoff from land, glacial and sea ice melt, and a varying presence of ArW in the study area.  
560 Bottom water salinity of about 35 is constant throughout the year, such that the seasonal cycle of  
561 density near the seafloor is controlled by temperature. Surface and bottom water temperatures  
562 rise towards summer, following atmospheric temperatures, regardless of which water mass is  
563 present in the area (Fig. 9). Temperatures increase towards summer from 2.5 to 6°C at the  
564 surface, and from 1.5 to 4°C at the bottom. The maximum temperature in the surface water is  
565 observed in July-August and one month later near the bottom. Winter surface and bottom  
566 temperatures vary between 1.5 and 2.5°C indicating that the water column is cooled down by  
567 heat loss to the atmosphere or surrounding waters (Nilsen et al., 2016). Warming of the water  
568 column in the study area throughout the year occurs through intermittent heat exchange with the

569 WSC that floods the shallow shelf (Nilsen et al., 2016), likely in a form of baroclinic eddies,  
 570 which are abundant in this region (Appen et al., 2016, Hattermann et al., 2016).



571

572 **Figure 8.** TS diagrams based on S800-model data for the study area, monthly average over the  
 573 period July 2005 – July 2010 (a); monthly average in May, June and July as these months are  
 574 when the surveys were conducted (b). Colours indicate month of the year as shown in the legend.  
 575 Background contour lines show isopycnals ( $\text{kg m}^{-3}$ ).



576

577 **Figure 9.** Annual cycle of bottom and surface seawater temperature and salinity in the study area,  
 578 modelled with S800-model. Lines show mean values for the study area, bars indicate spatial  
 579 variability.

580

### 581 3.6 Eddy driven seasonal dispersion on the shelf

582  
583 Our observations indicated a large spatial variability of dissolved methane concentrations,  
584 alongside limited vertical penetration of dissolved methane from the sources at the seafloor  
585 towards the sea surface irrespective of vertical stratification. Based on this, we propose that  
586 lateral advection near the seafloor plays an important role in dispersing methane horizontally  
587 away from the seep locations. The continuous replacement of methane enriched water with water  
588 containing low methane concentrations allows efficient dissolution of methane released in  
589 bubbles from the sediments.

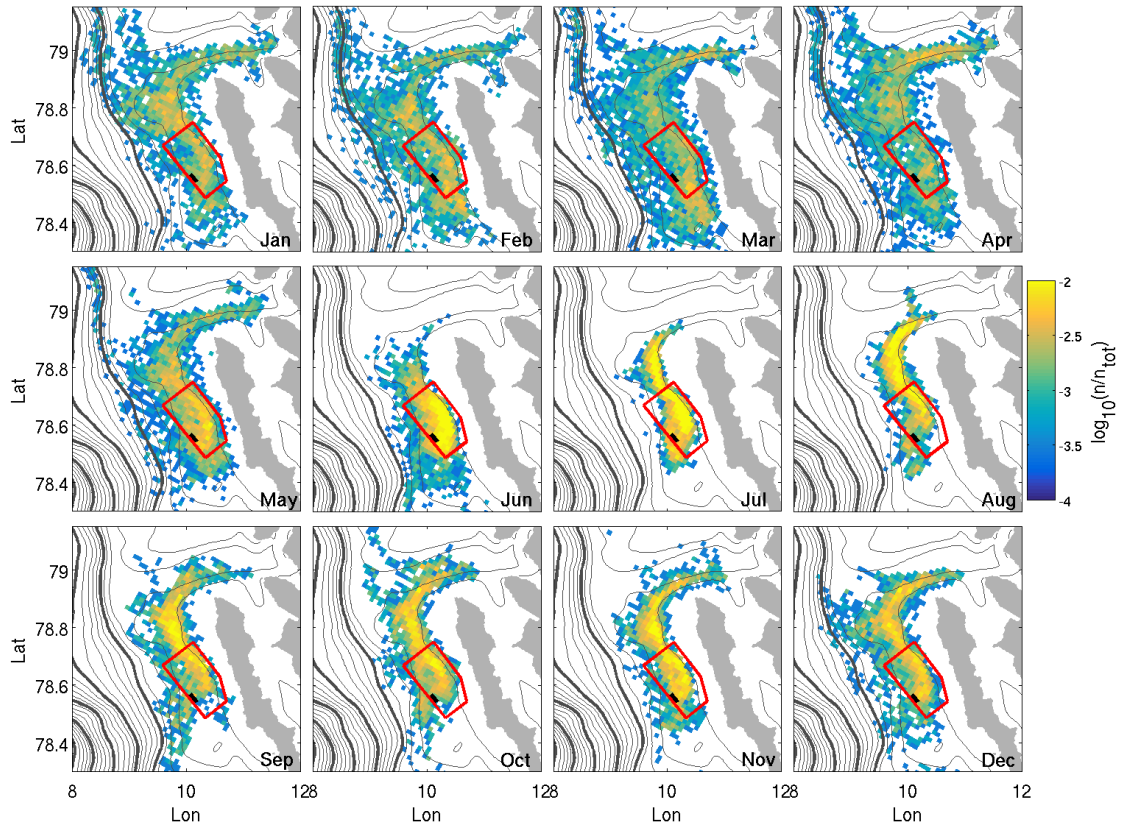
590 As shown on Fig. 6 (g-i), dissolved methane was spread horizontally in the bottom layer  
591 during June-14 and July-15 while it was more concentrated around the source in May-16. The  
592 water mass analysis suggests that this variability in horizontal dispersion is related to different  
593 circulation patterns on the shelf. As previously discussed, circulation of waters on the shallow  
594 shelf west of PKF is influenced by the combination of the WSC and superimposed local factors  
595 and their seasonality. The sole presence of the AW on the shelf in June-14 for example, which led  
596 to high dispersion of dissolved methane above the bottom, can be explained by an AW flooding  
597 event from the WSC over the shelf (Nilsen et al., 2008; Nilsen et al., 2016). Thereby, the lateral  
598 transport of waters above the PKF shelf during such flooding events disperses the dissolved  
599 methane and reduces the residence time of dissolved methane above gas flares.

600 While the WSC core generally flows further offshore than the shallow PKF shelf  
601 (Aagaard et al., 1987), instabilities of the WSC result in formation of numerous eddies that  
602 transport AW onto the shallow shelf (Appen et al., 2016, Hattermann et al. 2016, Wekerle et al.,  
603 2017). The transport occurs across the slope near the seafloor and plays an important role in the  
604 exchange of AW with shelf waters in our study region (Tverberg and Nøst, 2009). We propose  
605 that the observed large dispersion of dissolved methane above the bottom during the June-14  
606 survey is a result of eddy activity on the shallow shelf, and that eddies play an important role in  
607 the cross-frontal transport of waters and its constituents.

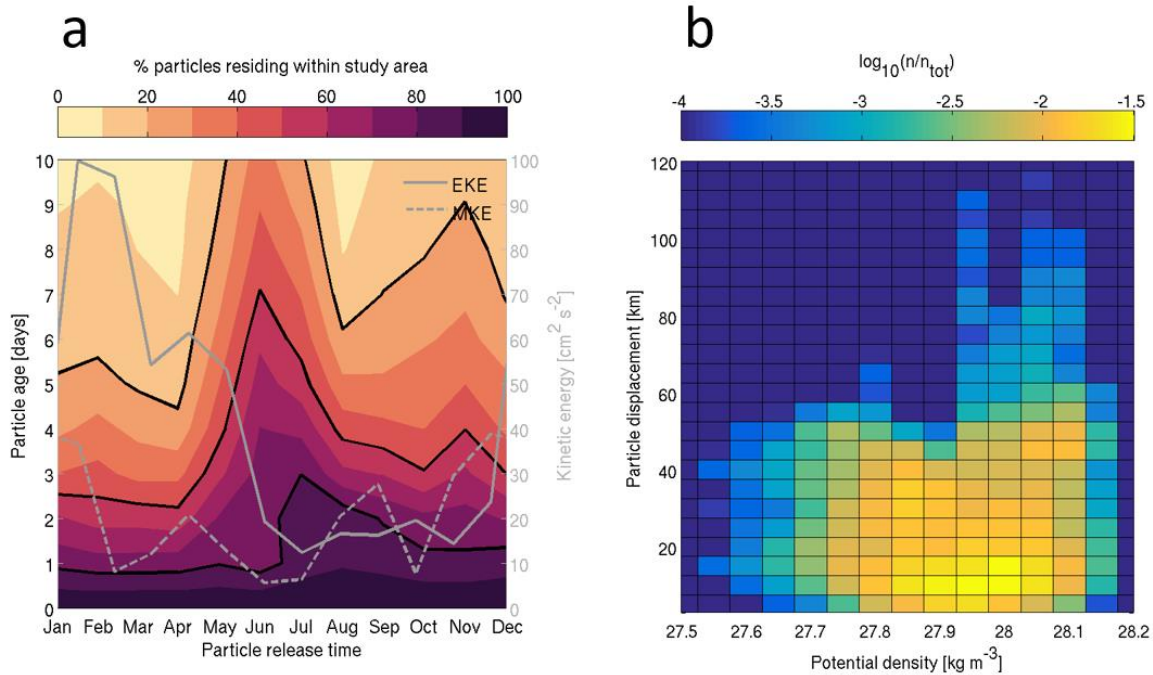
608 Appen et al. (2016) found increased eddy kinetic energy (EKE) and enhanced baroclinic  
609 instability in the WSC in winter and spring and it is likely that this seasonality will affect the  
610 number of flooding events over the shallow shelf and the residence time of methane above gas  
611 flares. To investigate the relationship between the seasonality of eddy activity and the variability

612 of dissolved methane dispersion on the shelf, we used the S800-model to run numerical  
613 experiments releasing and tracking particles simulating methane in our most intense flare area  
614 (see Methods 2.7). The particles are freely advected by the 3D model velocity field and provide a  
615 first order assessment of the role of the circulation in methane dispersion. The buoyancy driven  
616 motion of bubbles and the aerobic oxidation of dissolved methane will add further complexity to  
617 the dispersion process, but as discussed in section 3.3, are likely to be of secondary importance  
618 compared to the advective controls. In the numerical experiment, we observed a clear seasonality  
619 in particle dispersion with a much wider area being covered by the particles from January to May  
620 as opposed to a limited area of high particle concentrations during the summer months (Fig. 10).

621 During all months, the particles are mainly advected northward along the shelf and into  
622 the Kongsfjorden Trough that crosses the shelf at 79° N. However, in winter and spring, the  
623 pattern becomes more dispersive and particles are advected westward off the shelf, suggesting a  
624 greater influence of the WSC on water mass exchange with the shallow shelf area. The residence  
625 time within our study area follows the seasonal evolution of EKE (Fig 11a), with 50 % (80 %) of  
626 the released particles having left the study area after 3 days (6 days) between January to April,  
627 when EKE in the study area is largest. Furthermore, particles with the largest displacement (up to  
628 80 – 100 km five days after the release, Fig. 11b) are associated with the highest seawater density  
629 of 27.9 – 28.1 kg m<sup>-3</sup>, which is consistent with the hypothesis that methane is efficiently  
630 dispersed by eddies that lift dense AW onto the shelf (Tverberg and Nøst, 2009, Hattermann et al.  
631 2016). Although our observations during a 3-day period in each year do not resolve the seasonal  
632 cycle seen in the model, they support this principal mechanism, with the most dispersed methane  
633 concentrations being observed during the June-14 and July-15 surveys when AW was present in  
634 the bottom layer. Thus, our combination of observations and modelling suggests that eddies play  
635 an important role in dispersing outgassing methane over the continental shelf and in controlling  
636 the water column methane content, with potential direct implications for methane related  
637 biogeochemical processes.



638  
 639 **Figure 10.** Monthly maps of particle dispersion 5 days after the particle release between 80 and  
 640 100 m water depth at the positions indicated by the black rectangle. Colours indicate the number  
 641 of particles per grid cell normalized by the total number of particles in the respective month,  
 642 using a logarithmic scale. The red polygon delineates the location of the sampling sites, contours  
 643 show the isobaths with 100 m intervals thicker lines indicating 500 m intervals.



645

646 **Figure 11.** (a) Time series showing the residence time of particles within the study area indicated  
 647 by the red polygon in Fig. 10 (colour shade), together with monthly averaged mean (MKE) and  
 648 eddy (EKE) kinetic energy (right axis), averaged for the same region. Black curves indicate times  
 649 when 20 %, 50 % and 80 % of particles have left the study area. (b) Two-dimensional histogram  
 650 of particle displacement vs. potential density at the particle position after five days. Colours  
 651 indicate the normalized frequency of occurrence on a logarithmic scale, showing that many  
 652 particles remain within 20 km of the source and that the largest displacements are associated with  
 653 the highest densities.

654

#### 655 4 Conclusions

656 To our knowledge, this is the first study of the water column above cold methane seeps  
 657 that combines a multiyear series of oceanographic surveys with stations positioned on a grid  
 658 within a defined polygon. This study clearly benefits from the grid station design when compared  
 659 to more frequently conducted single synoptic transects. Acquiring data in a 4D array in time and



660 space allowed us to evaluate the methane content in the entire water body above methane flares  
661 and clearly identify the major processes mediating water column methane content and transport.

662 Our results suggest the possibility of enhanced methane flux from the sediments triggered  
663 by elevated bottom water temperature in the absence of underlying gas hydrate. In light of  
664 warming waters of the Arctic Ocean, not only gas hydrate containing sediments, but all methane  
665 gas bearing sediments could potentially become sources of methane release into the water  
666 column. Further study of the processes involved and the links between gas bearing sediments and  
667 bottom water temperature is required to improve our understanding.

668 Comparison between the three different hydrographic regimes observed across the three  
669 surveys reveals that most of the released methane in our shallow shelf area remains in the bottom  
670 and intermediate waters irrespective of the strength of stratification. Therefore, hypotheses by  
671 e.g. Schneider von Deimling et al., (2011), who suggested that all methane could be liberated to  
672 the atmosphere from shallow shelf areas as a result of a well-mixed water column and absence of  
673 stratification appear not to be valid in our shallow shelf study area. Small amounts of methane  
674 could be liberated to the atmosphere, but mainly as a result of strong winds increasing the rate of  
675 air-sea gas exchange, not weak stratification.

676 As expected, we find the horizontal advection to be the main mechanism controlling the  
677 dispersion of dissolved methane on the Prins Karls Forland shelf instead of vertical transport. In  
678 particular, our results highlight the role of mesoscale eddies in controlling the methane content  
679 above, dispersion around, and displacement away from gas flares. This implies that eddies and  
680 horizontal dispersion may also have important effects on methane related biogeochemical process  
681 and the magnitude of different methane sinks. For example, one could anticipate that a potential  
682 for methane sink through MOx could be higher when eddy activity is high in winter and spring  
683 season, because by dispersing dissolved methane over a larger area, eddies promote delivery of  
684 dissolved methane to methane oxidizing bacteria that consequently capture and consume this  
685 methane. Further seasonal measurements and/or process oriented modelling will be required to  
686 scrutinize these ideas, but these results could considerably shift our understanding of the  
687 seasonality of sinks of dissolved methane and allow better estimates of the balance between

688 amounts of methane released from sediments, methane liberated into the atmosphere, and  
689 methane removed from the system through microbial processes.

### 690 **Declaration of competing interest**

691 The authors declare that they have no known competing financial interests or personal  
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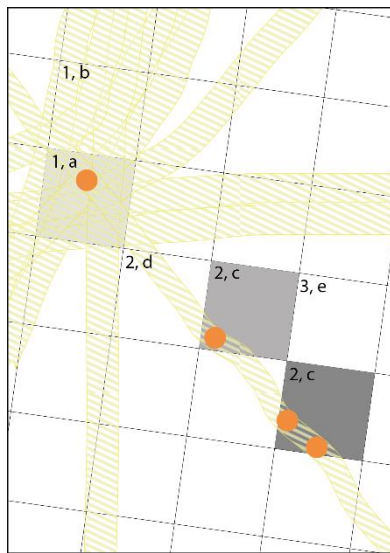
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971

972 **Appendices**

973 **Appendix A. Method for scaling up the flow rates**

974  
975 Because the fraction of the study area covered by the echosounder was small and slightly different  
976 between the three surveys, we applied a scaling up procedure including a Kriging interpolation to facilitate  
977 comparison between surveys (Fig. A1). The entire area was gridded into cells of  $100 \times 100$  m, and, as a  
978 result three types of cells were considered: 1) completely covered, 2) partly covered or 3) not covered by  
979 the echosounder beam footprint. For each cell, we estimated methane flow rates: (a) If one or several  
980 flares were detected within cell type 1, the estimated flow rate was applied for the entire cell area. (b) In  
981 the absence of flares in cell type 1, the flow rate was set to zero. (c) If one or several flares were within  
982 cell type 2, the sum of the flow rates within the cell was normalized by the fraction of the cell covered by  
983 the beam footprint. (d) In the absence of flares in cell type 2, the flow rate was set to zero. (e) For cell type  
984 3, (no data acquired) we interpolated flow rates from neighbouring cells. In order to find a smooth and  
985 plausible flowrate distribution, a  $3 \times 3$  low-pass filter and the Kriging interpolation method embedded in  
986 ArcGIS was applied. Finally, to calculate the mean flux ( $\text{mol m}^{-2} \text{s}^{-1}$ ) in the entire area, the sum of the  
987 scaled up flow rates were normalized by the survey area (Table 1).  
988



989  
990 **Figure A1.** Schematic of flowrate interpolation. Squares indicate 100x100 meter grid cells where the  
991 darkness indicates the relative summed flow rates within each cell. Yellow-hashed areas indicate the  
992 echosounder beam coverage and dots indicate flares. Cell types 1 – 3 and interpolation schemes a – e are  
993 described in the text.

994

995 **Appendix B. Calculation of methane content in the water column**

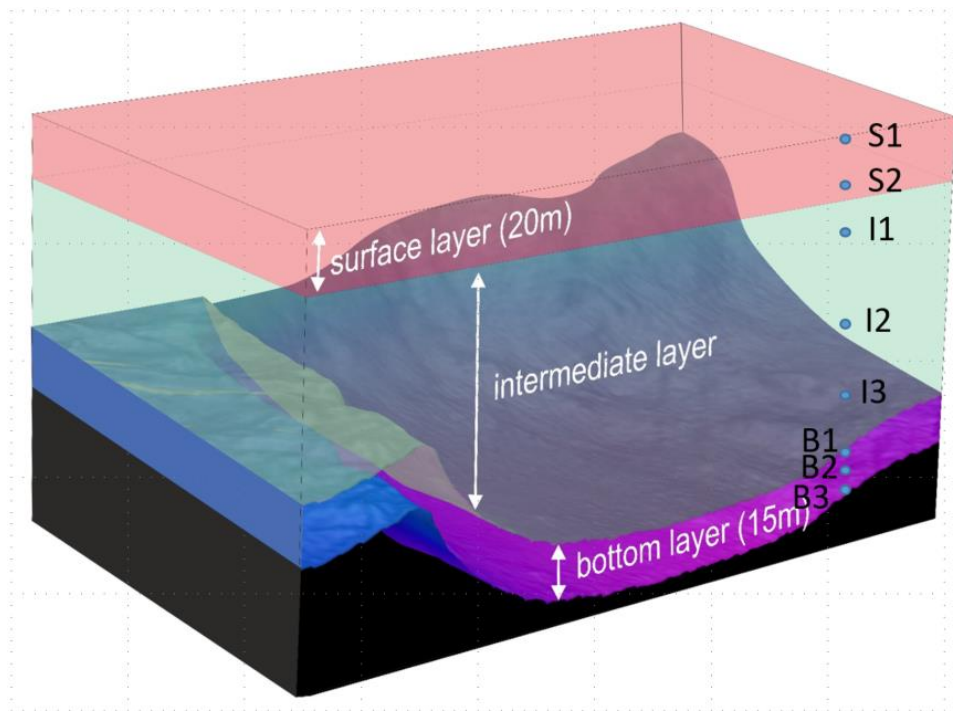
996  
997 Methane content in the water column was calculated by integrating distinct methane concentration values  
998 over depth. For this approach, we approximated the definite integral linearly by applying the trapezoid  
999 rule. For each discrete sampling point shown on Fig. B1, we had a corresponding depth ( $Z$ , m) and  $\text{CH}_4$   
1000 concentration ( $C$ ,  $\text{nmol L}^{-1}$ ). To determine the vertically integrated methane content ( $C_{\text{int}}$ ) for every depth  
1001 (meter) of water column between sampling points ( $\text{nmol L}^{-1} \text{ m}$ ) we interpolated linearly as follows:

$$1002 \quad C_{\text{int}} = (C_{S1} + C_{S2}) / 2 * (Z_{S2} - Z_{S1}) \quad (\text{Eq. B1})$$

1003  
1004 We then summed all  $C_{\text{int}}$  in each layer and multiplied by  $10^3$  to obtain methane content per  $\text{m}^2$  in every  
1005 layer ( $\text{nmol m}^{-2}$ ) for each of the CTD stations.

1006  
1007 To account for spatial sampling irregularity between CTD stations, we determined the area-weighted  
1008 average of the  $\text{CH}_4$  content for each layer. For this, we created a grid between longitudes  $9.5^\circ \text{ E}$  and  $10.8^\circ$   
1009  $\text{E}$  and latitudes  $78.4^\circ \text{ N}$  and  $78.7^\circ \text{ N}$  with bin sizes of  $0.01 \times 0.01^\circ$  in both directions. The resulting grid  
1010 included  $101 \times 201$  points. We then projected  $C_{\text{int}}$  for each layer and station onto this grid using the Matlab  
1011 function `griddata` for horizontal interpolation. Finally, we calculated the area-weighted average using the  
1012 Matlab function `mean2` of the gridded data. Then we scaled up (multiplied) the area-weighted averages for  
1013 each layer to the size of the investigated area ( $423 \text{ km}^2$ ), yielding the total methane content (in mol) for  
1014 each layer.

1015  
1016



1017  
 1018 **Figure B1.** Schematic shows the bottom layer (0-15 meters above seafloor), the intermediate layer (15  
 1019 meters above seafloor to 20 m water depth) and the surface layer (20 m water depth to sea surface). The  
 1020 blue dots show discrete sampling points in the surface (S1, S2), intermediate (I1, I2, I3) and bottom (B1,  
 1021 B2, B3) layer.  
 1022

1023 **Appendix C. Calculations of the sea-air methane flux**

1024 The sea-air methane flux  $F$  ( $\text{mol m}^{-2} \text{s}^{-1}$ ) was calculated according to Wanninkhof et al.  
1025 (2009):

1026 
$$F = k(C_w - C_o), \text{ (Eq. C1)}$$

1027 where  $k$  is the gas transfer velocity ( $\text{m s}^{-1}$ ),  $C_o$  is the methane concentration ( $\text{mol m}^{-3}$ ) at the  
1028 ocean surface in presumed equilibrium with the atmosphere and  $C_w$  is the measured  
1029 concentration of methane ( $\text{mol m}^{-3}$ ) in the well-mixed surface layer, typically measured at 5 m  
1030 water depth. The flux is positive and the ocean emits methane into the atmosphere if the  
1031 measured concentration in the surface layer is greater than the equilibrium concentration.  $C_o$   
1032 ( $\text{mol m}^{-3}$ ) is defined as:

1033 
$$C_o = \beta pCH_4, \text{ (Eq. C2)}$$

1034 where  $\beta$  is the Bunsen solubility ( $\text{mol m}^{-3} \text{atm}^{-1}$ ) of methane in seawater (Wiesenburg and  
1035 Guinasso, 1979):

1036 
$$\beta = \exp[-68.8862 + 101.4956 \left(\frac{100}{T_w}\right) + 28.7314 \left[\ln\left(\frac{T_w}{100}\right)\right] + S(-0.076146 +$$
  
1037 
$$0.04397 \left(\frac{T_w}{100}\right) - 0.0068672 \left(\frac{T_w}{100}\right)^2], \text{ (Eq. C3)}$$

1038 where  $T_w$  is the water temperature (K) and  $S$  is the salinity.

1039  $pCH_4$  is the partial pressure of methane in the air, derived from the mixing ratio of methane in the  
1040 atmosphere  $xCH_4$  ( $\text{mol mol}^{-1}$ ) measured by the on board CRDS at a height of 22.4 m (1902 ppb in  
1041 June-14, 1917 ppb in July-15 and 1955 ppb in May-16). The  $pCH_4$  was calculated according to  
1042 Pierrot et al., (2009):

1043 
$$pCH_4 = xCH_4 * [P_{atm} - P_{wvapor}], \text{ (Eq. C4)}$$

1044 accounting for the atmospheric pressure  $P_{atm}$  (atm) measured by the meteorological station on  
1045 board, and the water vapor pressure  $P_{wvapor}$  (atm) calculated according to Weiss and Price (1980):

1046 
$$P_{wvapor} = \exp[24.4543 - 67.4509 \left(\frac{100}{T_A}\right) - 4.8489 \ln\left(\frac{T_A}{100}\right) - 0.000544S], \text{ (Eq. C5)}$$

1047 where  $T_A$  is the air temperature (K) from the ships' meteorological station and  $S$  is the salinity of  
1048 spray in overlaying atmosphere, here assumed equal to the salinity of surface water.

1049 The gas transfer velocity  $k$  is wind dependent and calculated as described in Graves et al. (2015)  
1050 and references therein:

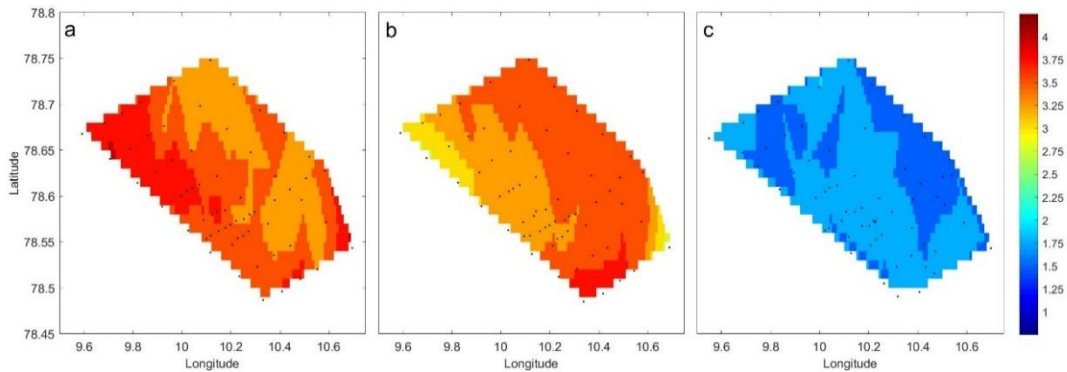
1051 
$$k = 0.24 * u_{10}^2 \left(\frac{Sc}{660}\right)^{-0.5}, \text{ (Eq. C6)}$$

1052 where  $u_{10}$  ( $\text{m s}^{-1}$ ) is the wind speed at 10 m above the sea surface, recalculated from the wind  
1053 speed  $u_{meas}$  ( $\text{m s}^{-1}$ ) measured by the ships' anemometer at height 22.4 m ( $z_{meas}$ ) after Hsu et al.,  
1054 1994:

1055 
$$u_{10} = u_{meas} * \left(\frac{z_{meas}}{10}\right)^{-0.11}, \text{ (Eq. C7)}$$

1056 The Schmidt number  $Sc$  in Eq. C6 is the non-dimensional ratio of gas diffusivity and water  
1057 kinematic viscosity, and was defined as 677 in accordance with Wanninkhof et al., (2009).

1058 **Appendix D. Additional information on bottom water temperature**



1059  
1060  
1061 **Figure D1.** Bottom water temperature during the June-14 (a), July-15 (b) and May-16 (c) surveys.  
1062

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